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Neutron diffraction study of the boride $Y(Fe_{11.04}Ti_{0.52}B_{0.43})Ti_{0.39}$: a new phase of a ThMn₁₂-type intermetallic interstitial compound

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Abstract. A sample with nominal composition YFe_{10.5}TiB_{0.5} has been prepared through careful processing. X-ray diffraction analysis revealed that the main phase in the sample is in the tetragonal ThMn₁₂-type structure (space group, 14/mmm), and a small amount (less than 3%) of α -Fe exists as a second phase. Neutron diffraction measurements indicated that, in the tetragonal structure, the boron atoms occupy the 8i sites, substituting for a few Ti atoms, while the substituted Ti atoms enter the interstitial 2b sites. The real composition of the main phase was determined to be Y(Fe_{11.04}Ti_{0.52}B_{0.43})Ti_{0.39}.

1. Introduction

The structure and magnetic properties of ternary compounds of the type $RFe_{12-x}M_x$ ($R \equiv$ rare earth; $M \equiv Ti$, V, Cr, Mo, W, etc) have been studied by numerous investigators [1-4]. These compounds crystallize in the tetragonal $ThMn_{12}$ structure in which there is only one crystallographic site for the rare earth and three non-equivalent sites for the M and Fe atoms. High Curie temperatures and high uniaxial anisotropies near room temperature make these systems promising starting materials for permanent-magnetic application [5, 6].

It has been discovered that the rare-earth-iron intermetallics of the $RFe_{12-x}Ti_x$ type can absorb moderate quantities of nitrogen [7]. Neutron diffraction measurements not only confirm that the nitrides maintain the tetragonal structure of their virginal compounds $RFe_{12-x}Ti_x$ but also indicate that the nitrogen atoms occupy the interstitial 2b sites, which leads to drastic changes in magnetic properties [7]. To date, the main disadvantage of the nitrided rare-earth-iron intermetallics is their thermal instability, which means they are easy to convert to RN_x and α -Fe phases at high temperatures.

In our previous work [8], we prepared various borides with the nominal composition $RFe_{12-x}(Ti,V)_xB_y$, trying to investigate the effects of boron atoms in the $ThMn_{12}$ -type rare-earth-iron intermetallics. Magnetic measurements, including high-field magnetization measurements and AC initial-susceptibility measurements, have shown that all the Curie temperatures, the iron moments and the anisotropies of the compounds undergo some changes upon the addition of boron [8]. Therefore, the position of boron atoms is the key to understanding the structure of the boride and interpreting the effects of boron on crystal-field and magnetic properties. In order to determine the positions of boron in these borides, we investigated the structure of $YFe_{10.5}TiB_{0.5}$ alloy by means of x-ray diffraction and neutron diffraction.



Figure 1. A part of the x-ray diffraction pattern of the (YFe_{10.5}TiB_{0.5}) sample.

2. Experimental details

The sample was prepared by arc melting the elemental constituents in a purified argon atmosphere, with the nominal composition $YFe_{10.5}TiB_{0.5}$; the melting point is 1286 °C (determined by the differential thermal analysis). To ensure homogeneity, the ingot was inverted and remelted. Then, the ingot wrapped in Mo foil was sealed in an evacuated quartz tube filled with high-purity argon and was annealed at 950 °C for 2 weeks, followed by rapid cooling to room temperature. The ThMn₁₂-type structure of the main phase was identified using x-ray diffraction with Cu K α radiation.

Neutron diffraction measurements were performed on the powder diffractometer at the heavy-water research reactor at the Institute of Atomic Energy at Beijing. The powder sample was packed in a thin-wall cylindrical vanadium cell of 0.5 cm diameter and 5 cm high. The diffraction data were collected at room temperature with a neutron wavelength of 1.184×10^{-8} cm by scanning from $2\Theta = 10^{\circ}$ to 80° in steps of 0.15° . The diffraction patterns obtained were analysed by means of Rietveld's [9] profile technique. For the coherent scattering lengths, we used the values 0.7750×10^{-12} cm for yttrium, 0.9540×10^{-12} cm for iron, -0.3300×10^{-12} cm for titanium and 0.5300×10^{-12} cm for boron. Indication of the agreement between observed and calculated intensities, both nuclear and magnetic, is given by the discrepancy factor R_{total} :

$$R_{\text{total}} = \sum_{i} \left| S_i(\text{obs}) - \frac{1}{c} S_i(\text{calc}) \right| / \sum_{i} S_i(\text{obs})$$



Figure 2. A part of the neutron diffraction pattern of the (YFe_{10.5}TiB_{0.5}) sample.

where $S_i(obs)$ and $S_i(calc)$ are the observed and calculated integrated intensities, to be summed over all reflections, and c is a scaling factor. Similar expressions are used to obtain the separate R_{nucl} - and R_{magn} -values.

3. Results and discussion

The results of x-ray diffraction with Cu K α radiation are shown in figure 1. It is confirmed from the pattern that the main phase in the sample is of the ThMn₁₂-type structure, and there is a little α -Fe.

The neutron diffraction diagram obtained at room temperature is shown in figure 2. The observed diffraction lines are interpreted by nuclear and magnetic coherent scattering. The diffraction patterns of the sample can be indexed in a polyphase program, while the main phase is in a body-centred tetragonal cell (space group, I4/mmm), with parameters $a = (8.509 \pm 0.004) \times 10^{-8}$ cm and $c = (4.808 \pm 0.003) \times 10^{-8}$ cm.

The observed and calculated intensities for the sample are listed in table 1.

The results of the refinement procedure for the main phase, including the coordinates of the crystallographic sites, the occupancy factors for different atoms in these sites and the magnetic moments are given in table 2.

The position parameters for Y, Fe and Ti are very close to those obtained for $YFe_{11}Ti$ [10]. The Fe atoms are found to occupy the 8i, 8j and 8f sites, whereas the Ti atoms reside in the 8i sites. The most important result is the location of boron atoms, as shown in table 1.

Table 1. Calculated and observed intensities at 300 K.

| hkl | 20 (deg) | $I_{calc}(nucl)$ | Icalc(mag) | Icalc(total) | Iobs(total) |
|------------|----------|------------------|------------|--------------|-------------|
| 200 | 15.996 | 694 | 48 | 742 | 742 |
| 101 | 16.261 | 212 | 10 | 222 | 183 |
| 301 | 28.048 | 910 | 9 | 919 | 953 |
| 002 | 28.514 | 377 | 0 | 377 | 505 |
| 400 | 32.315 | 2935 | 50 | 2985 | 3046 |
| 321 | 32.453 | 3884 | 43 | 3927 | 3993 |
| 202 | 32.862 | 2867 | 11 | 2878 | 2916 |
| 420 | 36.254 | 1919 | 40 | 1959 | 2030 |
| 411 | 36.379 | 1393 | 34 | 1427 | 1429 |
| 222 | 36.749 | 3934 | 28 | 3962 | 3851 |
| 440 | 46.351 | 490 | 3 | 493 | 431 |
| 521 | 46.453 | 115 | 0 | 115 | 103 |
| 422 | 46.757 | 411 | 1 | 412 | 394 |
| 303 | 50.212 | 275 | 0 | 275 | 283 |
| 512 | 51.169 | 238 | 0 | 238 | 195 |
| 620 | 52.208 | 186 | I | 187 | 161 |
| 611 | 52.301 | 185 | 0 | 185 | 160 |
| 323 | 53.042 | 1469 | 2 | 1471 | 1297 |
| 541 | 55.055 | 588 | 4 | 592 | 652 |
| 442 | 55.324 | 247 | 1 | 248 | 273 |
| 413 | 55.770 | 373 | 2 | 393 | 654 |
| 532 | 50.002 | 455 | 4 | 430 | 484 |
| 6031 | 57.719 | 1310 | 4 | 1320 | 1955 |
| 710 | 58 026 | 272 | 4 | 280 | 205 |
| 550 | 58 036 | 235 | 1 | 236 | 233 |
| 004 | 59.016 | 876 | ů L | 876 | 926 |
| 640 | 60 222 | 247 | 1 | 748 | 253 |
| 622 | 60.561 | 1535 | 4 | 1539 | 1530 |
| 721 | 62,830 | 367 | Ó | 367 | 433 |
| 730 | 63.988 | 117 | 0 | 117 | 186 |
| 712 | 66.753 | 437 | 1 | 438 | 452 |
| 552 | 66.753 | 372 | 0 | 372 | 384 |
| 800 | 67.637 | 888 | 1 | 889 | 898 |
| 651 | 67.716 | 120 | 0 | 120 | 121 |
| 642 | 67.955 | 925 | 1 | 926 | 952 |
| 613 | 68.351 | 112 | 0 | 112 | 118 |
| 404 | 68.905 | 1355 | 0 | 1355 | 1468 |
| 820 | 70.016 | 344 | I | 345 | 368 |
| 543 | 70.720 | 367 | 1 | 368 | 375 |
| 424 | 71.265 | 1034 | 1 | 1035 | 1013 |
| 732 | 71.503 | 191 | 0 | 191 | 183 |
| 660 | 72.360 | 184 | 0 | 184 | 165 |
| 633 | 73.054 | 975 | 1 | 976 | 870 |
| 531 | 74,730 | 4/1 | 0 | 4/1 | 440 |
| 214 | 74.743 | 602 | 0 | 407 | 607 |
| 040 877 | 70.902 | 433 | 0 | 433 | 407 |
| 723 | 77 641 | 255 | 0 | 255 | 245 |
| 444 | 78 168 | 406 | 0 | 406 | 379 |
| 534 | 79,299 | 244 | õ | 244 | 228 |
| 662 | 79.529 | 178 | õ | 178 | 166 |
| 761 | 81.554 | 320 | õ | 320 | 302 |
| 930 | 82.600 | 169 | ã | 169 | 165 |
| 624 | 82.669 | 160 | õ | 160 | 157 |
| 325 | 83.337 | 631 | 0 | 631 | 658 |

Table 2. Coordinates of sites, occupancy factors for different atoms and magnetic moments in different sites. $a = (8.509 \pm 0.004) \times 10^{-8}$ cm; $c = (4.808 \pm 0.003) \times 10^{-8}$ cm. $R_{\text{nucl}} = 6.48\%$; $R_{\text{magn}} = 8.46\%$; $R_{\text{total}} = 5.51\%$.

| Atom | Site | x | у | z | Number of atoms | μ _z (μ _B) |
|------|------------|----------|----------|----------|-----------------|-------------------------------------|
| Y | 2a | 0.000 00 | 0.000 00 | 0.000 00 | 2.00 | 0.00 |
| Fe | 8 i | 0.35610 | 0.000 00 | 0.000 00 | 6.09 | 1.83 |
| Ti | 8i | 0.35610 | 0.000 00 | 0.000 00 | 1.05 | 0.00 |
| В | 8i | 0.35610 | 0.000 00 | 0.000 00 | 0.87 | 0.00 |
| Fe | 8j | 0.277 00 | 0.500 00 | 0.000 00 | 8.00 | 1.71 |
| Fe | 8f | 0.250 00 | 0.250 00 | 0.250 00 | 8.00 | 1.66 |
| Ti | 2b | 0.000 00 | 0.000 00 | 0.500 00 | 0.78 | 0.00 |



Figure 3. Schematic representation of the structure of $Y(Fe_{11.04}Ti_{0.52}B_{0.43})Ti_{0.39}$.

The boron atoms occupy the 8i sites, substituting for some of the Ti atoms. Surprisingly, it is found that the substituted Ti atoms can enter the interstitial position 2b, like nitrogen atoms. Thus the real formula composition of the main phase is $Y(Fe_{11.04}Ti_{0.52}B_{0.43})Ti_{0.39}$.

A schematic representation of the structure of $Y(Fe_{11.04}Ti_{0.52}B_{0.43})Ti_{0.39}$ is shown in figure 3.

From the analysis of our neutron diffraction data, it is believed that the B atoms take the substitution position and substitute for the transition-metal atoms which enter the interstitial position. The α -Fe content in the sample is about 3%.

We have examined in the same way another sample with the nominal composition $(YFe_{11}Ti_{0.8}B_{0.2})$ [11]. The results shows that the main phase also has the ThMn₁₂-type structure, and the boron atoms prefer to occupy the substitution position, with the substituted Ti atoms in the 2b site. There is still a little α -Fe in the sample (less than 4%).

In conclusion, from the analysis of our neutron diffraction data, it can be derived that the boron atoms in the ThMn₁₂-type intermetallics prefer to occupy the substitution position, totally different from the nitrogen atoms which occupy only the interstitial position. Meanwhile, it is found that the substituted transition-metal atoms can enter the interstitial position 2b. Therefore, this work will be helpful for understanding the structure of borides of ThMn₁₂ type and its relationship to the magnetic properties.

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